Effect of Metal Ring Setting Outside HTSC Bulk Disk on Trapped Field and Temperature Rise in Pulse Field Magnetizing

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Abstract—In order to enhance the trapped field in cryo-cooled HTSC bulks using pulse field magnetizing (PFM), a metal ring (stainless steel 304 and/or Al) has been tightly set onto the Sm-BaCuO bulk disk and the relation between the total trapped flux Φ_T^P, the trapped field B_T^P, and the temperature rise ΔT, has been investigated as a function of the applied pulse field B_{ex}. The Φ_T^P and B_T^P values are enhanced about $10 \sim 20\%$ by the metal ring due to the reduction in the temperature rise ΔT. These results suggest that a part of the generated heat Q due to the flux motion in the peripheral region promptly transfers to the metal ring and the heat transfer to the cold stage is improved by the ring setting.

Index Terms—High T_c bulk superconductors, metal ring setting, pinning and viscous loss, pulse field magnetizing, temperature measurement.

I. INTRODUCTION

IN VIEW of the practical applications of high-T_c bulk superconductors (HTSCs) as a high strength bulk magnet for a magnetic levitation system and so on, pulse field magnetizing (PFM) as well as the static field-cooled magnetizing (FCM) has been intensively investigated because of the relatively compact and inexpensive setup. The field trapped by PFM is, however, lower than that attained by FCM at temperatures below 77 K. The main reason has been attributed to the large heat generation due to the dynamic motion of the magnetic flux against the vortex pinning force F_p and the viscous force F_v. We have studied the temperature rise ΔT and trapped field B_T^P on the surface of cryo-cooled YBaCuO and SmBaCuO bulks during PFM and pointed out the importance of the ΔT reduction for the B_T^P enhancement [1]–[4]. The total generated heat Q was estimated using the specific heat C of the bulk and the maximum ΔT. It was pointed out that the B_T^P values as functions of the initial temperature T_s and applied pulse field B_{ex} can be understood on the basis of the trapped field (B_T^{FC}) by FCM vs.

the temperature (T) diagram, i.e., by the decrease of the critical current density J_c associated with the temperature rise. In order to enhance the B_T^P and Φ_T^P values by PFM, the reduction of the ΔT is an indispensable issue. The iteratively magnetizing pulsed-field method with reducing amplitude (IMRA) [5], locating yoke pieces around a bulk [6], a use of vortex-type coils [7] and a multi-pulse technique with step-wise cooling (MPSC) [8] have been attempted to suppress the heat generation during PFM. A B_T^{FC} of 17.24 T at 29 K in YBaCuO bulk has been realized by FCM, enhancing the thermal conduction of the bulk by impregnating a high thermal conductivity alloy into drilled holes [9].

In this study, we set a metal ring (Al and/or stainless steel 304) tightly on a SmBaCuO bulk disk. The following effects are expected from the metal ring set on to the HTSC bulk disk. First, since the heat generation due to PFM occurs mainly in the peripheral region, the metal ring attached to the bulk disk periphery can easily and promptly receive the heat. Second, since the ab-plane thermal conductivity κ_{ab} is far larger than the c-axis κ_c, the generated heat can easily reach the metal ring along the radial paths of the disk, then transferring to the cold stage through the highly conductive ring. These two effects should contribute to the ΔT reduction of the bulk. Finally, the metal ring setting helps to mechanically reinforce the bulk disk [10].

II. EXPERIMENTAL PROCEDURE

A highly c-axis oriented SmBaCuO bulk superconductor (Dowa Mining Co., Ltd) with 45 mm diameter and 15 mm thickness was used. This consisted of four growth sector regions (GSRs), divided by the growth sector boundaries (GSBs) [2]. The bulk is composed of $SmBa_2Cu_3O_y$ (Sm123) and Sm_2BaCuO_5 (Sm211) with the molar ratio of Sm123 : Sm211 $= 1.0 : 0.3$, 15.0 wt.% Ag_2O powder, and 0.5 wt.% Pt powder. The bulk was uniformly impregnated by epoxy resin in vacuum and then the epoxy resin layer on the surface of the bulk disk was removed. The stainless steel (SUS304) or the aluminum (Al) ring with 4 mm in thickness and 15 mm in height, was fixed onto the bulk disk using apiezon-N grease. We call the ring-attached sample as SUS-Sm or Al-Sm. Since a large heat generation due to the eddy current took place in the Al ring ($\Delta T > 6$ K at 100 K), the ring was cut with a gap of 1 mm as shown in Fig. 1. Resultantly, ΔT was diminished to within 0.2 K. Since the Al ring did not contact tightly with the bulk due to the gap, another SUS ring
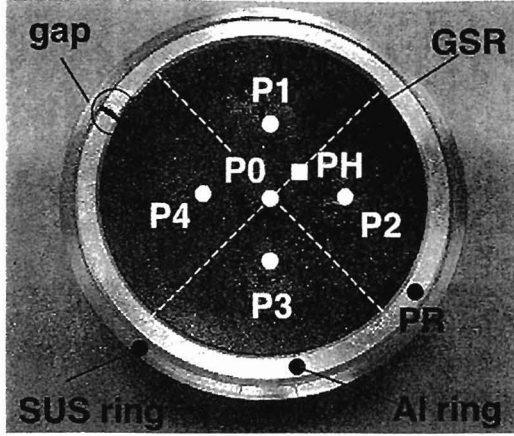
Fig. 1. Photograph of the SmBaCuO-bulk disk fitted with the W(Al+SUS)-ring. The positions of the temperature and magnetic field measurements (P0 ~ P4, PR and PH) are shown.

TABLE I
SUMMARY OF SIZES AND MATERIAL PARAMETERS OF Sm-BULK AND RINGS

	Sm-bulk	SUS-ring	Al-ring	W-ring (Al-SUS)
I.D. (mm ϕ)	-	45	45	45-53
O.D. (mm ϕ)	45	53	53	53-55
Height (mm)	15	15	15	15-15
V (mm^3)	23.84	9.24	9.24	9.24-2.54
C_V (J/cm^{3}K) at 40 K	0.48	0.38	0.23	-
C (J/K) at 40 K	11.44	3.47	2.08	3.04
C ratio	1.0	0.30	0.18	0.27
κ at 40K (mW/cmK)	130 (ab) 30 (c)	55	~5000	-

with 1 mm in thickness was set outside the Al-ring. Hereafter, we denote this sample with the Al+SUS (W) ring the W-Sm. Table I summarizes the sizes, the volume of each metal ring, and the specific heat C_V (J/cm^{3}K), the heat capacity C (J/K) at 40 K of the Sm-bulk and each metal ring, and the ratio of the heat capacity of the metal ring to that of the Sm-bulk. The heat capacity increases by 30%, 18% and 27% for the SUS-ring, Al-ring and W-ring, respectively, relative to that for the bare Sm bulk. The ab-plane and c-axis thermal conductivity κ at 40 K of the Sm-bulk and that of SUS and Al are also shown [11]. C_V and κ were measured for the present specimens [12].

The bulk disk was tightly stacked on the sapphire plate (45 mm in diameter and 20 mm in thickness) attached to the cold stage of a helium refrigerator. The initial stage temperature T_s was kept at 40 K. The temperatures, $T0$ at the center of the bulk (P0) and, $T1 \sim T4$ at P1 ~ P4 were monitored by fine chromel-constantan thermocouples adhered to the upper bulk surface by GE7031 varnish. P1 ~ P4 were situated on the central radial lines of each GSR by 9 mm apart from P0. The temperature of the metal ring TR was also measured at PR. The total trapped magnetic flux Φ_T^P and the distribution of the trapped
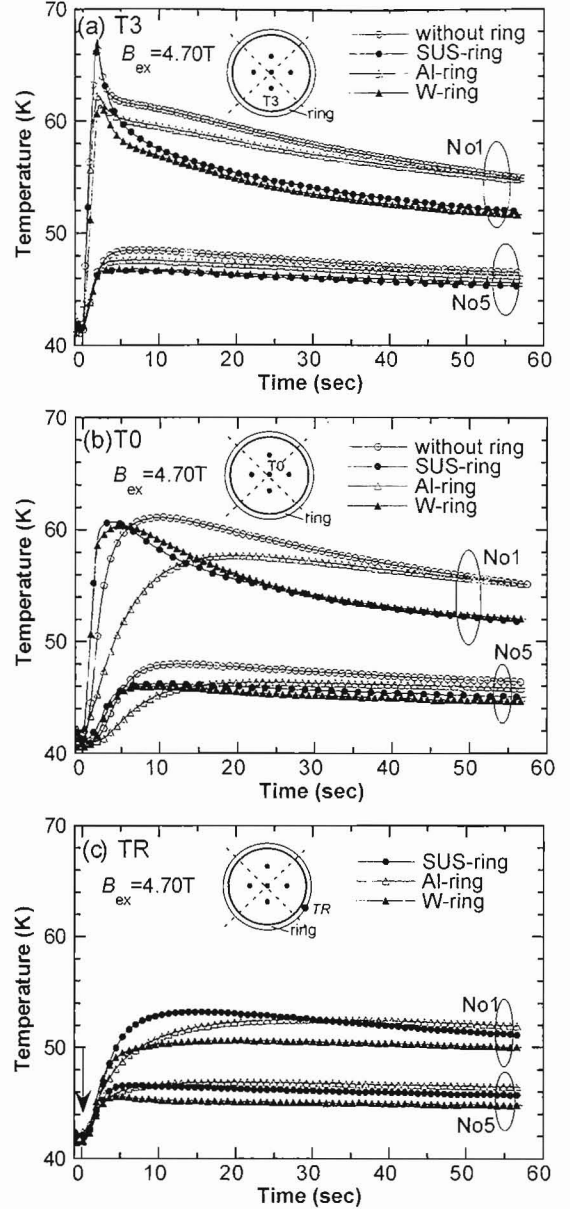

Fig. 2. Examples of the time evolution of temperatures $T(t)$ of the Sm-bulk without ring, with SUS-ring, with Al-ring and with W(SUS+Al)-ring at (a) $T3$, (b) $T0$ and (c) TR after applying the No1 and No5 pulse fields of $B_{ex} = 4.70$ T.

magnetic flux density $B_T^{3\ \text{mm}}$ were measured using an axial type Hall sensor (F.W. Bell, model BHA 921), which scanned 3 mm above the bulk surface stepwise with a pitch of 1.2 mm. The trapped field B_T^P on the bulk surface was measured by the Hall sensor adhered to the position at PH with a 2.5 mm distance from P0. Five iterative magnetic pulses (No1 ~ No5) with the same amplitude B_{ex} from 3.83 T to 6.04 T (rise time: 12 ms) were applied sequentially after re-cooling to T_s. $T(t)$, Φ_T^P and B_T^P were measured at each stage.

III. RESULTS AND DISCUSSION

Fig. 2(a) shows the time dependence of temperature $T3(t)$ at P3 after applying the No1 and No5 pulse of $B_{ex} = 4.70$ T. For the bare Sm-bulk (B-Sm) without the ring, the magnetic
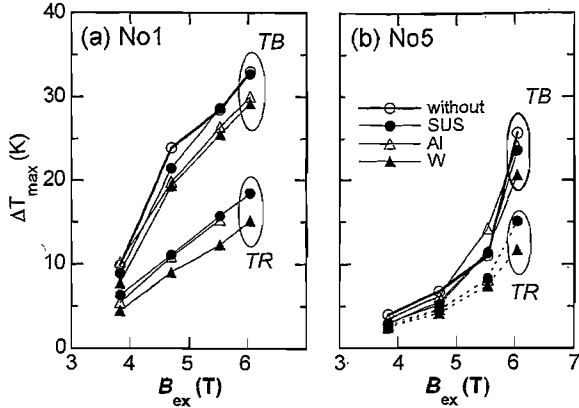
Fig. 3.   The maximum temperature rises ΔTB_{max} and ΔTR_{max} at the bulk and the ring for (a) the No1 and (b) No5 pulses as a function of B_{ex}. The ΔTB_{max} values are the averaged ones over $\Delta T0_{max}$ to $\Delta T4_{max}$.
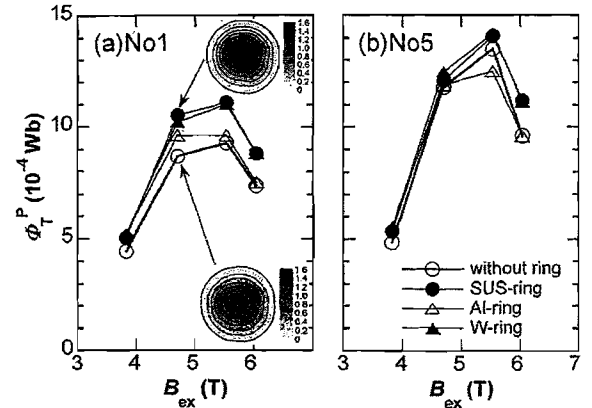

Fig. 4.   The total trapped flux Φ_T^P for each metal ring setting after (a) the No1 and (b) the No5 pulse application, as a function of B_{ex}. The typical conical trapped field distributions $B_T^{3\ mm}$ are displayed for the SUS-ring setting and bare Sm-bulk.

flux preferentially intrudes into the bulk from the GSRs containing P2, P3 and P4 and the sharp and large temperature rises ($\Delta T = 25$ K) take place within 3 s at these positions. $T(t)$ recovers to the initial temperature T_s after 15~20 min. $T3(t)$ of SUS-Sm decreases faster after the $T3(t)$ maximum, although the maximum temperature rise is similar to that of B-Sm. This result means that a part of the generated heat Q promptly transfers to the SUS-ring. The maximum ΔT value at $T3$ decreases to ~20 K for Al-Sm and W-Sm. The rate of temperature decrease $(-dT(t)/dt)$ after the peak is small for Al-Sm but that for the W-Sm is large and almost the same as the SUS-Sm. As for the iterative pulse field applications, ΔT is the largest for the No1 pulse and decreases for the succeeding pulses [2]. The increment of the trapped field ΔB_T^P is also the largest for the No1 pulse, followed by a gradual increase for the No2 and No3 pulses. For the No5 pulse application, the $T(t)$ peak disappears for all the settings. ΔT is slightly smaller for the SUS-Sm and W-Sm than for the bare bulk and Al-ring setting.

Fig. 2(b) presents the $T0(t)$ at P0 for each metal ring setting. $T0(t)$ rises up latest among $T0 \sim T4$ because the distance from the major heat source is usually the longest [2]. For the SUS-Sm and W-Sm, $T0(t)$ reaches a maximum faster than that for B-Sm. These results mean that the flux motion is enhanced in the central region of the bulk by the SUS-Sm and W-Sm and the magnetic fluxes are accumulated faster in the bulk center. The anomalously long time constant for the temperature rise for the Al-ring setting may result from the high thermal contact resistance between the bulk and Al-ring.

Fig. 2(c) shows the $TR(t)$ for each metal ring setting. The maximum ΔT values for SUS-Sm and W-Sm are 12 K and 10 K, respectively, and the time constant for the temperature rise of $TR(t)$ is longer for the SUS-Sm and W-Sm than that of $T3(t)$. The temperature rise in the metal ring is mainly due to the heat conduction from the bulk and the temperature rise in the Al-ring is slower because of the high thermal contact resistance.

Figs. 3(a) and (b) summarize the maximum temperature rise of the bulk ΔTB_{max} and the ring ΔTR_{max} after the No1 and No5 pulse as a function of B_{ex}. ΔTB_{max}, the averaged value at $\Delta T0_{max} \sim \Delta T4_{max}$, increases with increasing B_{ex} for both No1 and No5 pulses. ΔTB_{max} of B-Sm is the largest and is reduced by the metal ring setting. Especially, the reduction of

the ΔTB_{max} is the largest for W-Sm and Al-Sm; e.g., from 23 K to 18 K for the No1 pulse of $B_{ex} = 4.70$ T. ΔTR_{max} also increases with increasing B_{ex} for both No1 and No5 pulses and is about 50 to 80% of ΔTB_{max}. ΔTR_{max} is the largest for the SUS-ring and the smallest for the W-ring.

Figs. 4(a) and 4(b) present the total trapped flux Φ_T^P after the No1 and No5 pulse as a function of B_{ex}. For the No1 pulse, Φ_T^P of B-Sm is smallest for each B_{ex} and the Φ_T^P value is enhanced by the metal ring setting. Especially, for the SUS-Sm and W-Sm, the Φ_T^P values after the 4.70 T and 5.53 T pulses increase about 20% compared with those of the bulk without ring. The increase may mainly come from the decrease of the temperature rise due to the ring setting, which results in the increase of the effective critical current density. The Φ_T^P enhancement is observed for Al-Sm, which suggests that the good thermal contact between the ring and the Sm-bulk is of vital importance. The trapped field B_T^P measured at PH is also enhanced by the metal ring setting; for example, $B_T^P (= 2.95$ T) for W-Sm is about 10% larger than that of B-Sm ($= 2.71$ T) after the No5 pulse of $B_{ex} = 5.53$ T. The typical trapped field distributions $B_T^{3\ mm}$ are shown in the inset of Fig. 4(a) in which a conical field distribution can be confirmed.

Neglecting the heat drained to the cold stage, the generated heat Q by the PFM operation can be estimated using the following equation,

$$Q = \int_{T_s}^{T_s + \Delta TB_{max}} (C_{Bulk} V_{Bulk}) dT$$

$$+ \int_{T_s}^{T_s + \Delta TR_{max}} (C_{Ring} + V_{Ring}) dT, \quad (1)$$

where C_{Bulk} and C_{Ring} are the specific heat (J/cm^3K), and V_{Bulk} and V_{Ring} are the volume of the Sm-bulk and the metal ring, respectively. Fig. 5 shows the estimated Q values as a function of B_{ex} for each metal ring setting. For the No1 pulse, the Q value of the bulk with each metal ring is smaller than that of B-Sm. This result suggests that the heat transferred from the bulk is partially drained to the cold stage through the ring. For
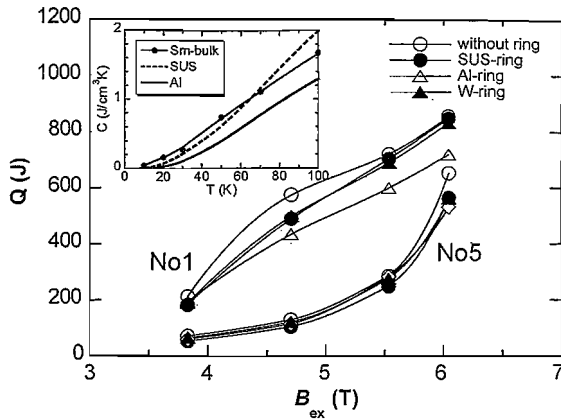
Fig. 5. The estimated generated heat Q as a function of B_{ex} for each metal ring setting. The inset indicates the heat capacity C of the Sm-bulk, SUS304 and Al.

the No5 pulse, The Q value of B-Sm is nearly the same as the total Q of the bulk and the metal ring, irrespective of the species and volume of the metal ring. This result means that the heat is generated mainly in the bulk and is transferred to the metal ring nearly adiabatically for the No5 pulse. The heat flow to the cold stage may remain small in case $\Delta T B_{max}$ is not so large.

IV. SUMMARY

In order to enhance the trapped field in cryo-cooled HTSC bulks by the pulse field magnetizing (PFM) method, the effect of a metal ring set around the SmBaCuO bulk disk has been investigated by measuring the total trapped flux Φ_T^P, the trapped field B_T^P and the temperature rise ΔT for the successive application of pulse fields B_{ex} with various strengths. Important experimental results and conclusions obtained in this study are summarized as follows.

(1) The Φ_T^P and B_T^P values are enhanced about 10~20% by the metal ring setting compared with those of the bare bulk without the metal ring. The main origin is the reduction of ΔT of the bulk during PFM. Since the heat generation mostly takes place in the peripheral region of the bulk disk in PFM, the metal ring set around the bulk is more effective than the insertion of a high-thermal-conductive alloy on the bulk center.

(2) The SUS- and W(Al+SUS)-rings behave as a good heat reservoir which contribute to the ΔT reduction in the

bulk. The good thermal conductance between the bulk periphery and ring is necessary to reduce ΔT and to enhance the Φ_T^P and B_T^P values.

(3) The metal ring with a high thermal conductivity also enhances the thermal drain for the generated heat Q to transfer to the cold stage.

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REFERENCES

[1] H. Fujishiro, T. Oka, K. Yokoyama, and K. Noto, "Time evolution and spatial distribution of temperature in YBCO bulk superconductor after pulse field magnetizing," *Supercond. Sci. Technol.*, vol. 16, pp. 809–814, 2003.

[2] ———. "Temperature rise in Sm-based bulk superconductor after applying iterative pulse fields," *Supercond. Sci. Technol.*, vol. 17, pp. 51–57, 2004.

[3] H. Fujishiro, T. Oka, K. Yokoyama, M. Kaneyama, and K. Noto, "Flux motion studies by means of temperature measurement in magnetizing processes for HTSC bulks," *IEEE Appl. Supercond.*, vol. 14, no. 2, pp. 1054–1057, 2004.

[4] H. Fujishiro, K. Yokoyama, M. Kaneyama, T. Oka, and K. Noto, "Estimation of generated heat in pulse field magnetizing for SmBaCuO bulk superconductor," *Physica C*, vol. 412–414, pp. 646–650, 2004.

[5] U. Mizutani, T. Oka, Y. Itoh, Y. Yanagi, M. Yoshikawa, and H. Ikuta, "Pulse-field magnetization applied to high-T_c superconductors," *Appl. Supercond.*, vol. 6, pp. 235–246, 1998.

[6] H. Ikuta, H. Ishihara, Y. Yanagi, Y. Itoh, and U. Mizutani, "Extracting the utmost from the high performance of Sm-Ba-Cu-O bulk superconductors by pulse field magnetization," *Supercond. Sci. Tech.*, vol. 15, pp. 606–612, 2002.

[7] T. Ida, H. Matsuzaki, Y. Akita, M. Izumi, H. Sugimoto, Y. Hondou, Y. Kimura, N. Sakai, S. Nariki, I. Hirabayashi, M. Miki, M. Murakami, and M. Kitano, "Magnetization properties for Gd-Ba-Cu-O bulk superconductors with a couple of pulsed-field vortex-type coils," *Physica C*, vol. 412–414, pp. 638–645, 2004.

[8] M. Sander, U. Sutter, R. Koch, and M. Kläser, "Pulsed magnetization of HTS bulk parts at T < 77 K," *Supercond. Sci. Tech.*, vol. 13, pp. 841–845, 2000.

[9] M. Tomita and M. Murakami, "High-temperature superconductor bulk magnets that can trap magnetic fields of over 17 tesla at 29 K," *Nature*, vol. 421, pp. 517–520, 2003.

[10] M. Morita et al., "Trapped field and mechanical properties of QMG magnet," in *Proc. 1998 Int. Workshop Superconductivity*, Tokyo, 1998, pp. 115–118.

[11] H. Fujishiro and S. Kohayashi, "Thermal conductivity, thermal diffusivity and thermoelectric power in Sm-based bulk superconductors," *IEEE Trans. Appl. Supercond.*, vol. 12, pp. 1124–1127, 2002.

[12] M. Ikebe, H. Fujishiro, T. Naito, and K. Noto, "Simultaneous measurement of thermal diffusivity and conductivity applied to Bi-2223 ceramic superconductors," *J. Phys. Soc. Jpn.*, vol. 63, pp. 3107–3114, 1994.